

REFERENCES

- (1) MTW CHAPTER 21
- (2) YORK, "KINEMATICS & DYNAMICS OF GR" IN "SOURCES OF GRAV. RADIATION" (L. STARR ed)
- (3) WALD, APP. E.2 ; CHAPTER 10
- (4) ARNOWITT, DESER ; MISNER (1962) "THE DYNAMICS OF GR" IN "GRAVITATION: AN INTRODUCTION TO CURRENT RESEARCH" (L. WITTEN ed)

- ADM MOTIVATION WAS, AGAIN, PREPARATION FOR QUANTIZATION OF GR; FORMALISM TURNED OUT TO BE EXCELLENT BASIS FOR COMPUTATIONAL (NUMERICAL) ASSAULT ON EINSTEIN EQUATIONS

APPROACHES

- (1) "COORDINATE-FUL" (MTW): INTUITIVE, CONNECTS MORE DIRECTLY TO FORM OF E.O.M. USED IN PRACTICE
- (2) "COORDINATE-FREE" (YORK/WALD): PREFERABLE FOR DERIVATION OF E.O.M.

KEY POINT: MUST INTRODUCE A COORDINATE SYSTEM TO NUMERICALLY GENERATE A SPACE-TIME; I.E. CAN'T STAY COORDINATE FREE FOREVER

- WILL START WITH COORDINATE BASED APPROACH TO INTRODUCE CONCEPTS, THEN WILL GO OVER TO COORDINATE-FREE APPROACH TO DERIVE 3+1 EQNS

ULTIMATE GOAL: REFORMULATE

$$G_{ab} = E_{Ti} T_{ab}$$

AS SYSTEM OF FIRST-ORDER (IN TIME) PDE'S FOR THE GRAV. FIELD VBL'S, WHICH CAN THEN BE SOLVED AS AN "INITIAL-VALUE" OR "CAUCHY" PROBLEM

SHIFT IN PERSPECTIVE: UP TO NOW $G_{ab} = E_{Ti} T_{ab}$ DESCRIBED THE LINKAGE OF THE GEOMETRY OF SPACETIME (4-D) TO THE DISTRIBUTION OF MATTER-STRESS-ENERGY IN S.T.

NEW VIEW: GEOMETRY OF S.T. IS "TIME-HISTORY" (EVOLUTION) OF GEOMETRY OF A SPACELIKE HYPERSURFACE (3-D) => GEOMETRODYNAMICS; "VIEW" NOT UNIQUE SINCE THERE ARE CO'LY MANY WAYS OF "SLICING UP" GIVEN S.T. INTO A FAMILY OF S.L. HYPERSURFACES; NONE PREFERRED (I.E. PHYSICALLY MORE RELEVANT) IN GENERAL

SPLITTING SPACETIME INTO SPACE-PLUS-TIME (THE 3+1 SPLIT)

• SPACETIME IS 4-DIMENSIONAL MANIFOLD M , WITH LORENTZIAN SIGNATURE METRIC g_{ab} (-+++)

• INTRODUCE COORDINATES $\{x^a\} = \{t, x^i\}$ (MAY NOT COVER ENTIRE S.T., BUT WILL GENERALLY PRETEND THEY DO)

NOTATION / CONVENTIONS:

GREEK INDICES: (μ, ν etc.) 0, 1, 2, 3

"INTEGER" LATIN " (i, j, k, l, m, n) 1, 2, 3 (SPATIAL)

WILL ADOPT USUAL EINSTEIN SUMMATION CONVENTION FOR BOTH TYPES!

• DEMAND THAT Σ = COINST SURFACES (HYPERSURFACES),

$\Sigma(\pm)$, ARE SPACELIKE; I.E. IF DISTINCT EVENTS P, P'

HAVE COORDS (t, x^i) , $(t, x^{i'})$ THEN $ds^2(P, P') > 0$

• SAFEST TO VIEW t AS THE PARAMETER; MAY NOT (IN FACT, IN GENERAL WILL NOT) HAVE ANY PARTIC. SIGNIFICANCE AS A PHYSICAL TIME - SUCH A NOTION IS LARGELY MEANINGLESS IN GR

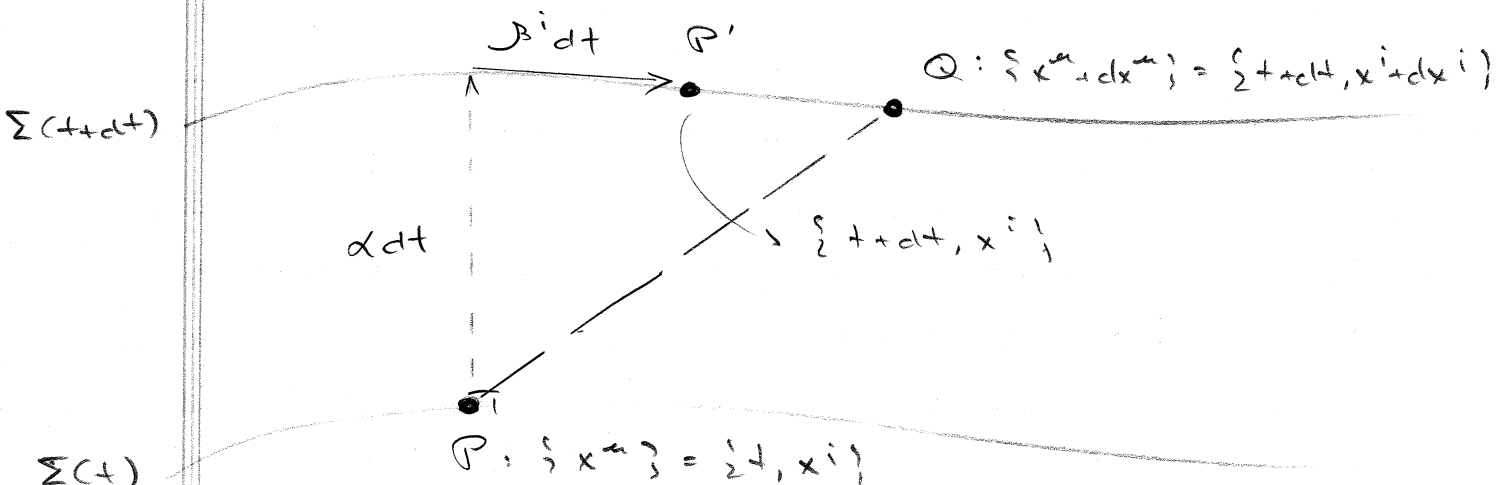
• EACH $\Sigma(t)$ IS A DIFFERENTIABLE MANIFOLD IN ITS OWN RIGHT, WITH A 3-METRIC ${}^{(3)}g_{ij} = {}^{(3)}g_{(ij)}$ WHICH IS INDUCED ON $\Sigma(t)$ BY THE 4-METRIC ${}^{(4)}g_{\mu\nu}$ OF THE ENVELOPING SPACETIME

• ANY GIVEN FOLIATION (= CHOICE OF TIME COORD = CHOICE OF SLICING) ALSO DEFINES A NATURAL, UNIT-NORM VECTOR FIELD, \vec{n} , WHICH IS NORMAL (ORTHOGONAL) TO THE SLICES, AND POINTS "TO THE FUTURE"

$$\vec{n} \cdot \vec{n} = {}^{(4)}g_{\mu\nu} n^\mu n^\nu = -1$$



SPACETIME DISPLACEMENT IN THE 3+1 SPLIT



"SPACETIME PITHACORAN TMM" ($\lim dt \rightarrow 0$) \Rightarrow

$$\text{distance}(P, Q)^2 = c^2 ds^2$$

$$= -c^2 dt^2 + {}^{(3)}g_{ij} (dx^i + \beta^i dt) (dx^j + \beta^j dt)$$

$$= (-c^2 + {}^{(3)}g_{ij} \beta^i \beta^j) dt^2 + 2 {}^{(3)}g_{ij} \beta^i dx^j dt$$

$$+ {}^{(3)}g_{ij} dx^i dx^j \quad (1)$$

NOTATION:

$\alpha \equiv \alpha(t, x^i)$: LAPSE FUNCTION ("LAPSE"): GIVES LAPSE OF PROPER TIME PER UNIT COORDINATE TIME FOR AN OBSERVER MOVING NORMAL TO THE SLICES

$\beta^j \equiv \beta^j(t, x^i)$: SHIFT VECTOR ("SHIFT"): 3-VECTOR DESCRIBING SHIFTING OF SPATIAL COORDINATES RELATIVE TO "NORMAL PROPAGATION"

TOGETHER $\{\alpha, \beta^i\}$ CONSTITUTE 4-FOLD COORD. FREEDOM of GR

DUAL VIEWS

(1) α, β^i ARE ESSENTIALLY FREELY SPECIFIABLE FIELDS (COORDINATE FREEDOM)

(2) SOME PRESCRIPTION FOR α, β^i MUST BE GIVEN "FROM OUTSIDE" - I.E. E.O.M. (CONSTRAINTS) ALONE WILL NOT, IN GENERAL, DETERMINE THEM ("GAUGE FIXING")

PHYS 327N THE 3+1 FORMULATION OF GR

(3)

NOTE: TENSORS SUCH AS β^i ARE DEFINED ON $\Sigma(t)$, AND ARE CALLED SPATIAL TENSORS

CLEARLY, ${}^{(3)}g_{ij}$ IS A SPATIAL TENSOR; IT HAS AN ASSOCIATED INVERSE ${}^{(3)}g^{ij}$ SATISFYING

$${}^{(3)}g^{ij} {}^{(3)}g_{jk} = \delta^i_k$$

INDICES OF 3-TENSORS ARE RAISED/LOWERED WITH ${}^{(3)}g^{ij}$, ${}^{(3)}g_{ij}$; THUS, WE CAN REWRITE (1) AS

$${}^{(4)}ds^2 = (-\alpha^2 + \beta^i \beta_i) dt^2 + 2\beta_j dx^j dt + {}^{(3)}g_{ij} dx^i dx^j$$

SO, WE HAVE

$${}^{(4)}g_{\mu\nu} = \begin{bmatrix} {}^{(4)}g_{00} & {}^{(4)}g_{0j} \\ {}^{(4)}g_{i0} & {}^{(4)}g_{ij} \end{bmatrix} = \begin{bmatrix} -\alpha^2 + \beta^k \beta_k & \beta_j \\ \beta_i & {}^{(3)}g_{ij} \end{bmatrix} \quad (2)$$

IN PARTICULAR, NOTE THAT

$${}^{(4)}g_{ij} = {}^{(3)}g_{ij}$$

I.E. SPATIAL COVARIANT COMPONENTS OF 4- AND 3-METRICS ARE IDENTICAL

THIS IS A GENERAL RESULT; GIVEN ANY 1-TENSOR OF TYPE $(0, k)$ (COVARIANT TENSOR), THE SPATIAL COMPONENTS OF THAT TENSOR CAN BE IDENTIFIED AS THE COMPONENTS OF A TYPE $(0, k)$ 3-TENSOR

WHY? RECALL THAT COVARIANT TENSOR COMPONENTS CAN BE

DEFINED IN TERMS OF THE ACTION OF THE TENSOR ON

THE COORDINATE BASIS VECTORS $\{^{(A)}\underline{e}_\mu\}$, $\mu=0, 1, 2, 3$

$$\text{E.G. } ^{(A)}t_{\mu\nu} = ^{(A)}t \left(^{(A)}\underline{e}_\mu, ^{(A)}\underline{e}_\nu \right)$$

$$\text{AND } ^{(A)}t_{ij} = ^{(A)}t \left(^{(A)}\underline{e}_i, ^{(A)}\underline{e}_j \right)$$

BUT CLEARLY THE $\{^{(A)}\underline{e}_i\}$ ARE PRECISELY THE
COORDINATE BASIS VECTORS $\{^{(B)}\underline{e}_i\}$ FOR $\Sigma(t)$ WITH
COORDINATES $\{x^i\}$; THIS INTERPRETING $^{(A)}t(\dots)$
AS A 3-TENSOR $^{(B)}t(\dots)$, WE NECESSARILY HAVE

$$^{(B)}t_{ij} = ^{(A)}t_{ij}$$

ON THE OTHER HAND, THE SPANAL MEMBERS $\{^{(A)}\underline{\omega}^i\}$ OF
THE DUAL BASIS $\{^{(A)}\underline{\omega}^a\}$, WILL NOT, IN GENERAL COINCIDE
WITH $\{^{(B)}\underline{\omega}^i\}$ - THE DUAL BASIS ON $\Sigma(t)$ DEPENDS ON
HOW $\Sigma(t)$ IS EMBEDDED IN THE S.T.

• THUS, IN GENERAL

$$^{(B)}t_{ij} \neq ^{(A)}t_{ij}$$

EXAMPLE: CONSIDER THE INVERSE 4-METRIC COMPONENTS
FROM (2)

$$^{(A)}g_{\mu\nu} = \begin{bmatrix} -1/\alpha^2 & \beta^j/\alpha^2 \\ \beta^i/\alpha^2 & ^{(B)}g_{ij} - \beta^i\beta^j/\alpha^2 \end{bmatrix} \quad (3)$$

(EXERCISE: VERIFY)

FROM (2) WE CAN ALSO COMPUTE THE USEFUL RESULT

$$\sqrt{-^{(4)}g} = \alpha \sqrt{{}^{(3)}g} \quad (A)$$

THE NORMAL VECTOR FIELD n^μ

EASIEST TO START WITH ASSOC. DUAL-VECTOR (ONE-FORM) FIELD, \underline{n}_μ

GEOM. INTERP of DUAL-VECTOR FIELD: LEVEL SURFACES of SCALAR FUNCTION \rightarrow DUAL NORMAL TO "INFINITESIMAL DISPLACEMENT" (VECTOR)



$$\langle \vec{v}, \underline{df} \rangle = v^\mu (df)_\mu = \# \text{ of level surfaces "pierced" by } \vec{v}$$

HERE, OUR SCALAR FUNCTION IS THE TIME COORDINATE t , WITH ASSOCIATED DUAL-VECTOR FIELD \underline{dt} , THEN

$$\underline{n} \propto \underline{dt}$$

OR IN COMPONENT FORM

$$n_\mu = (n_0, 0, 0, 0)$$

THEN, FROM

$$(A) \quad g^{\mu\nu} n_\mu n_\nu = -1$$

WE HAVE

SIGN CHOSEN SO THAT n^μ IS FUTURE-DIRECTED

$$n_\mu = (-\alpha, 0, 0, 0) \quad (5)$$

AND THEN

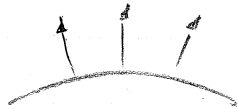
$$n^\mu = g^{\mu\nu} n_\nu = \left(\frac{1}{\alpha}, -\frac{\beta^i}{\alpha} \right) \quad (6)$$

EXTRINSIC CURVATURE

• THE INTRINSIC GEOMETRY of $\Sigma(t)$ IS DESCRIBED BY g_{ij} WHICH ENCODES ALL GEOM. INFO. WHICH MAY BE DEDUCED BY MAKING MEASUREMENTS ON $\Sigma(t)$ ALONE

• HOWEVER, A GIVEN 3-GEOMETRY (SLICE, HYPERSURFACE) MAY BE EMBEDDED IN SPACETIME IN INFINITELY MANY DISTINGUISHABLE (BY 4-D MEASUREMENTS) WAYS

EXAMPLE: 2D EMBEDDED IN 3D; FLAT SURFACE EMBEDDED WITH / WITHOUT EXTRINSIC CURVATURE



• THE MANNER IN WHICH $\Sigma(t)$ IS EMBEDDED CAN BE CHARACTERIZED BY INVESTIGATING THE CHANGE IN THE DIRECTION OF THE NORMAL FIELD AS A Fcn of POS ON $\Sigma(t)$ - THIS DEFINES THE EXTRINSIC CURVATURE TENSOR (ALSO THE SECOND FUNDAMENTAL FORM)

$$K_{ij} = -\nabla_i n_j = -\nabla_{(i} n_{j)} \quad (7)$$

ANY 3+1 THE 3+1 FORMULATION of GR

$$\nabla_i n_j = \partial_i n_j - \Gamma^{\mu}_{ij} n_{\mu}$$

BUT $n_{\mu} = (-\alpha, 0, 0, 0)$ SO

$$\begin{aligned} \nabla_i n_j &= \alpha^{(4)} \Gamma^0_{ij} \\ &= \alpha \left(g^{(4)00} \Gamma^0_{ij} + g^{(4)0c} \Gamma^c_{ij} \right) \end{aligned}$$

$$\Gamma^0_{ij} = \frac{1}{2} \left(g^{(4)}_{0i,j} + g^{(4)}_{0j,i} - \underline{g^{(4)}_{ij,0}} \right)$$

THUS

$$K_{ij} = -\frac{1}{2\alpha} \frac{\partial}{\partial t} g^{(3)}_{ij} + \dots$$

SO WE CAN VIEW THE EXTRINSIC CURVATURE AS THE "VELOCITY" OF THE 3-METRIC $g^{(3)}_{ij}$ ("CONJUGATE MOMENTA")